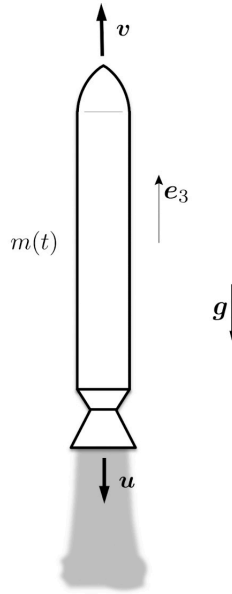


## Problem 1: Rocket launch

A rocket of initial mass  $m_i$  is launched from the earth at rest and accelerates until all of its fuel has been exhausted. The residual (final) mass of the rocket's structure is  $m_f$ . Assume that the fuel is burned at a constant positive rate  $\alpha$  and the exhaust gas is released vertically downward with a *relative* speed  $u = \text{const}$ . Neglect air resistance and assume the gravitational acceleration  $g$  (acting downwards) is constant.

Given:  $m_i, m_f, \alpha, u, g$



1. What is the velocity of the rocket  $v(t)$  at a given time  $t$  after launch?
2. What altitude  $x_3(t_f)$  has the rocket reached, when all its fuel has been burned at time  $t_f$ ?

1) Let's use the rocket equation in  $\underline{e}_3$ :

$$F_{\text{ext}} = m \frac{dv}{dt} + \frac{dm}{dt} \cdot (v - v_m) \quad (1)$$

where  $F_{\text{ext}} = -mg$  and  $(v - v_m) = u$  (relative)

- We can find  $m(t) = m_i - \alpha t$
- And the time at which all fuel is burned:

$$t_f = \frac{m_i - m_f}{\alpha}$$

- We can now rearrange (5) by dividing by  $m$  and multiplying with  $dt$

$$\Rightarrow dv + u \frac{dm}{m} = -g dt$$

now we can integrate:

$$\int_0^v dv + u \int_{m_i}^{m(t)} \frac{dm}{m} = -g \int_0^t dt$$

$$v + u \left[ \ln(m) \right]_{m_i}^{m(t)} = -gt$$

$$\int \frac{1}{x} dx = \ln(x)$$

$$\underline{\underline{v(t) = -gt - u \ln\left(\frac{m_i - \alpha t}{m_i}\right)}}$$

$$\ln(a) - \ln(b) = \ln\left(\frac{a}{b}\right)$$

- 2) We can integrate  $v(t)$  from  $t=0$  to  $t=t_f$

$$X_3(t_f) = -\frac{1}{2} g t_f^2 - u \int_0^{t_f} \ln\left(\frac{m_i - \alpha t}{m_i}\right) dt$$

- Now we substitute  $\mu = \frac{m_i - \alpha t}{m_i}$ ,  $\Rightarrow dt = -\frac{m_i}{\alpha} d\mu$

- Now we can integrate:  $\frac{m_f}{m_i}$  ← remember to also change the boundaries

$$u \int_0^{t_f} \ln\left(\frac{m_i - \alpha t}{m_i}\right) dt = u \int_1^{\frac{m_f}{m_i}} \ln(\mu) \cdot \frac{m_i}{\alpha} d\mu$$

$$\begin{aligned}
 \Rightarrow & u \cdot \frac{-m_i}{\alpha} \left[ \mu \ln(\mu) - \mu \right] \frac{m_f}{m_i} \\
 &= u \cdot \frac{-m_i}{\alpha} \left[ \frac{m_f}{m_i} \ln\left(\frac{m_f}{m_i}\right) - \frac{m_f}{m_i} + 1 \right] \\
 &= u \cdot \frac{1}{\alpha} \left[ \frac{m_f - m_i}{m_i} - \frac{m_f}{m_i} \ln\left(\frac{m_f}{m_i}\right) \right]
 \end{aligned}$$

$$\begin{aligned}
 \Rightarrow x_3(t_f) &= -\frac{1}{2} g \left( \frac{m_i - m_f}{\alpha} \right)^2 \\
 &\quad - \frac{u}{\alpha} \left[ \frac{m_f - m_i}{m_i} - \frac{m_f}{m_i} \ln\left(\frac{m_f}{m_i}\right) \right]
 \end{aligned}$$


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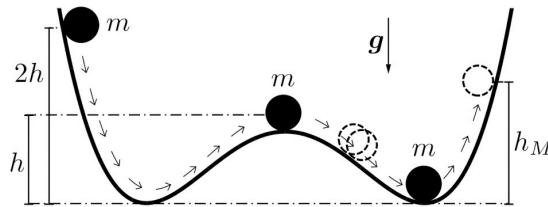


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## Problem 2: Chain reaction

Three particles of equal masses  $m$  and negligible dimensions are at rest on a frictionless curvilinear guide lying in the vertical plane, as shown below (the black particles). The first particle is released from rest from an initial height  $2h$ , and eventually impacts the second particle with a *perfectly plastic* collision. After the impact, these two particles slide together, until they collide with the third particle. The coefficient of restitution of this last impact is a known value  $e$ . What is the maximum height  $h_M$  reached by the last particle? Assume that the particles never detach from the guide and that gravity acts downwards, as shown.

Given:  $h, e, m, g$



We can use WEB and the collision equations to solve this problem.

- First we use WEB to find the velocity of the first particle at the first impact.

$$mgh = \frac{1}{2} m v_{1-}^2 \Rightarrow v_{1-} = \sqrt{2gh}$$

- Since the first impact is plastic ( $e=0$ ) the collision formula can be simplified to: ( $v_{1+} = v_{2+}$ , they stick!)

$$m v_{1-} + m v_{2-} = 2m v_{1+} \Rightarrow v_{1+} = v_{2+} = \frac{v_{1-}}{2}$$

- Now we can use WEB again to find  $v_1$  before the second collision:

$$\frac{1}{2} 2m v_{1+}^2 + 2mgh = \frac{1}{2} 2m v_{1-}^2$$

$$\Rightarrow v_{1-} = \sqrt{v_{1+}^2 + 2gh}$$

- The collision formula gives the velocity of the 3rd particle after the 2nd collision:

$$v_{3+} = \frac{(1+e) 2m}{m+2m} v_{1-}$$

- Now we can use WEB once again to find  $h_M$

$$mgh_M = \frac{1}{2} m v_{3+}^2$$

$$\Rightarrow h_M = \frac{v_{3+}^2}{2g} = \frac{5(1+e)^2}{g} h$$

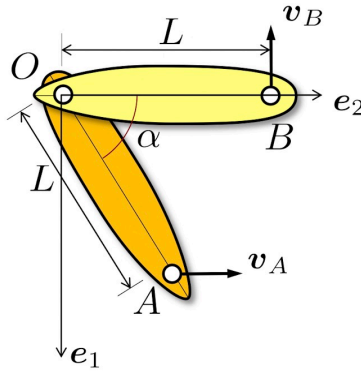
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### Problem 3: Kinematics of coupled rigid bodies

The system sketched below consists of two rigid bodies  $OA$  and  $OB$ , connected at point  $O$  by a joint that allows relative rotation. In the configuration shown, the two bodies enclose an angle  $\alpha$ . The velocities at points  $A$  and  $B$  are given by  $\mathbf{v}_A = v_A \mathbf{e}_2$  and  $\mathbf{v}_B = -v_B \mathbf{e}_1$ , respectively, and  $|\mathbf{r}_{AO}| = |\mathbf{r}_{BO}| = L$ . What is the magnitude  $\omega_{OA}$  of the angular velocity  $\boldsymbol{\omega}_{OA}$  of body  $OA$ ?

Given:  $v_A, v_B, L, \alpha$



We can use the velocity transfer formula to find  $\underline{v}_O$ , from both bodies:

$$\underline{v}_O = \underline{v}_A + \underline{\omega}_{OA} \times \underline{r}_{AO} = \underline{v}_B + \underline{\omega}_{OB} \times \underline{r}_{BO}$$

which we can write as:

$$v_A \underline{e}_2 + \omega_{OA} \underline{e}_3 \times (-L \sin(\alpha) \underline{e}_1 - L \cos(\alpha) \underline{e}_2) =$$
$$-v_B \underline{e}_1 + \omega_{OB} \underline{e}_3 \times (-L \underline{e}_2)$$

then we can carry out the vector products:

$$v_A \underline{e}_2 - \omega_{OA} L \sin(\alpha) \underline{e}_2 + \omega_{OA} L \cos(\alpha) \underline{e}_1 =$$
$$-v_B \underline{e}_1 + \omega_{OB} L \underline{e}_1$$

which, when we write it componentwise gives us two equations:

$$X_1: -V_B + \omega_{0B}L = \omega_{0A}L \cos(\alpha)$$

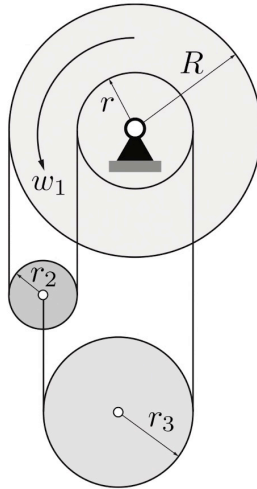
$$X_2: V_A - \omega_{0A}L \sin(\alpha) = 0$$

$$\Rightarrow \omega_{0A} = \frac{V_A}{L \sin(\alpha)}$$

### Problem 4: Pulley mechanism

As part of a pulley mechanism, three disks are connected to each other by non-slipping ropes, as shown below. What are the velocities of the centers of the two smaller rolls, if the large disk rotates with an angular velocity  $\omega_1$ , as shown? What are their angular velocities?

Given:  $r$ ,  $R$ ,  $\omega_1$



We introduce points A, B, C to more easily convert between the velocities:

- $V_A = R\omega_1$ ,  $V_B = r\omega_1$ ,  $V_C = -r\omega_1$  (1)

- As there is no slip, we know for the smallest roll

$$V_A = V_2 + r_2 \omega_2$$

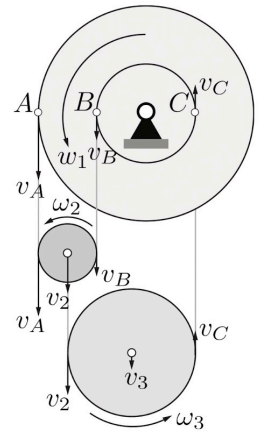
← This is essentially the transfer formula in 2D

$$V_B = V_2 - r_2 \omega_2 \quad (2)$$

- We can now use (1), (2) and  $r_2 = \frac{R-r}{2}$

to find  $\omega_2$ :

$$\omega_2 = \frac{V_A - V_2}{r_2} = \frac{R\omega_1 - \frac{1}{2}(R+r)\omega_1}{\frac{1}{2}(R+r)} = \omega_1$$



For the bottom roll we can do the same:

$$V_3 = V_2 - r_3 \omega_3, \quad V_C = V_3 - r_3 \omega_3$$

We also know  $r_3 = \frac{R+3r}{4}$ , so we can find  $\omega_3$

$$\omega_3 = \frac{V_2 - V_3}{r_3} = \frac{\frac{1}{2}(R+r)\omega_1 - \frac{1}{4}(R-r)\omega_1}{\frac{1}{4}(R+3r)} = \omega_1$$

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